

# Lecture 12 Electric Displacement

Note Title

11/3/2009

Ch. 4 HW Part 2: 15, 18, 19, 21, 28

In a dielectric, write  $\rho = \rho_b + \rho_f$

$$\text{Then } \nabla \cdot \vec{E} = \frac{\rho}{\epsilon_0} \iff \nabla \cdot \epsilon_0 \vec{E} = \rho_b + \rho_f$$

But  $\rho_b = -\nabla \cdot \vec{P}$ , so Gauss's law is

$$\boxed{\nabla \cdot \vec{D} = \rho_f}$$

$$\text{for } \boxed{\vec{D} = \epsilon_0 \vec{E} + \vec{P}}$$

$\vec{D}$  is the electric displacement.

$\hookrightarrow$  In integral form

$$\boxed{\oint \vec{D} \cdot d\vec{a} = Q_{\text{free}}}$$

$\oint \vec{D} \cdot d\vec{a} = Q_{\text{free}}$  is very useful if there is symmetry, for instance problem 15.

On the other hand,  $\nabla \times \vec{D} = \nabla \times \vec{P} \neq 0$

so  $\nabla \cdot \vec{D} = \rho_f$  isn't sufficient in finding  $\vec{D}$  if there isn't symmetry.

And  $D(\vec{r}) \neq \frac{1}{4\pi} \int \frac{\hat{r}}{r^2} \rho_f(\vec{r}') dV'$

o At an interface

$$D_{\text{above}}^{\perp} - D_{\text{below}}^{\perp} = \sigma_f$$

and

$$D_{\text{above}}^{\parallel} - D_{\text{below}}^{\parallel} = P_{\text{above}}^{\parallel} - P_{\text{below}}^{\parallel}$$

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The most common, and useful, application are linear dielectrics, where

$$\vec{P} = \epsilon_0 \chi_e \vec{E}$$

$\chi_e \leftarrow$  Electric susceptibility

$$\text{Then } \vec{D} = \epsilon_0 \vec{E} + \vec{P} = \epsilon_0 (1 + \chi_e) \vec{E}$$

$$\text{or } \vec{D} = \epsilon \vec{E} \quad \epsilon \leftarrow \text{permittivity}$$

$$\text{and define } \epsilon_r = \frac{\epsilon}{\epsilon_0} = 1 + \chi_e$$

$\uparrow$   
relative permittivity  
or dielectric constant.

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If a space is entirely filled w/ a homogeneous dielectric, then  $\vec{E}$  is reduced by  $\epsilon_r$ .

$\hookrightarrow$  or pretty much replace  $\epsilon_0$  w/  $\epsilon$ !

In particular, you can use  $\oint \vec{D} \cdot d\vec{a} = Q_{\text{free}}$  to solve capacitor problems, and  $C = \epsilon_r C_{\text{vac}}$ .

How about energy? To charge cap:  $W = \frac{1}{2} CV^2$   
 so if  $C \uparrow$ ,  $W \uparrow$

$$W = \frac{1}{2} \int \vec{D} \cdot \vec{E} d\tau \dots \text{energy in the fields.}$$

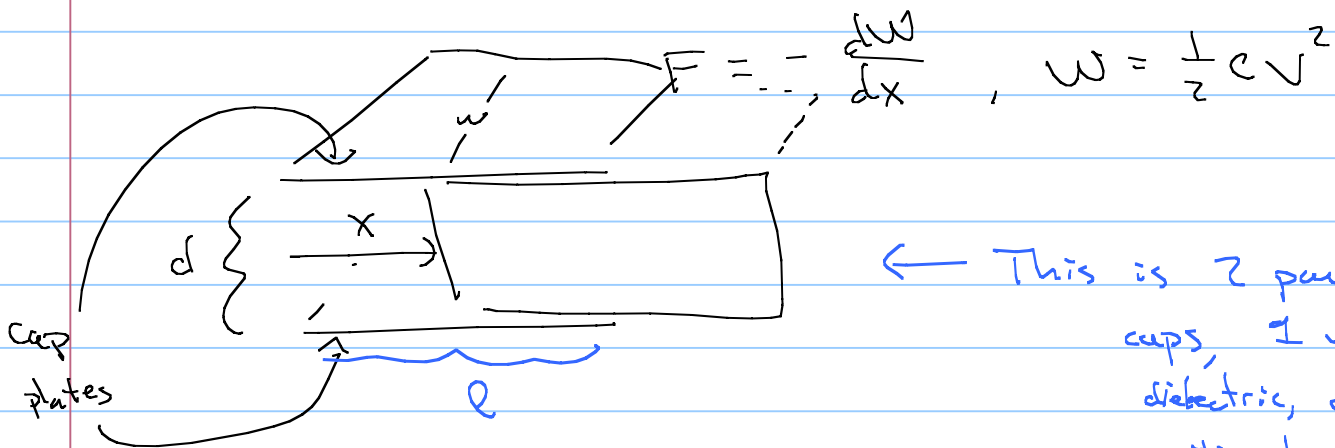
What about forces on dielectrics?

Dielectrics are pulled into the  $\vec{E}$  field.

- o Suppose you have a dielectric slab inside a capacitor and apply a force,  $F_{me}$ , to pull it out a distance  $dx$ .

↳ The work you do is  $dW = F_{me} dx$ ,

working against the electrical force on the slab



← This is 2 parallel caps, 1 w/ dielectric, one without

$$C = C_{\text{air}} + C_{\text{die}} = \frac{\epsilon_0 A_{\text{air}}}{d} + \frac{\epsilon_0 \epsilon_r A_{\text{die}}}{d}$$

$$C = \epsilon_0 \frac{xw}{d} + \frac{\epsilon_0 (1 + \chi_e)(l-x)w}{d}$$

$$C = \frac{\cancel{\epsilon_0 xw}}{d} + \frac{\epsilon_0 w l (1 + \chi_e)}{d} - \frac{\epsilon_0 (\cancel{x + \chi_e}) x w}{d}$$

$$\boxed{C = \frac{\epsilon_0 w}{d} (l \epsilon_r - \chi_e x)} \quad \text{eqn 4.62}$$

w/  $Q$  held fixed,  $Q = \sqrt{C} \Rightarrow W = \frac{Q^2}{2C}$

$$F = -\frac{dW}{dx} = \frac{1}{2} \frac{Q^2}{C^2} \frac{dC}{dx} = \frac{1}{2} \sqrt{2} \frac{dC}{dx}$$

$$\boxed{F = -\frac{\epsilon_0 \chi_e w}{2d} \sqrt{2}} \quad \dots \text{ck}$$