

So we've got

Schrödinger interpretation of Q.M.

Solve: (0)  $\hat{H}\Psi(x,t) = i\hbar \frac{\partial \Psi}{\partial t}$

if w/  $\Psi$  we have

(1)  $\int_N \Psi^* \Psi dx = 1$  solve for

Then  $\bar{\Psi}(x,t) \rightarrow N \Psi(x,t)$

= Normalized  
wavefunction

(2) Then  $\Psi^* \bar{\Psi} = \frac{\text{prob}}{\text{length}}$

$\Psi^* \Psi dx = \text{probability.}$

So what?

To do more.... We'll need to revisit  
Probability Theory.

no longer deterministic  
VIEW of particles!

Before doing that --- recall where we are  
at

Light = particle + wave  $\Rightarrow$  Wave Equat  $\rightarrow$  Wave Funct

$e^-$  = particle + wave  $\Rightarrow$  Wave Equat  $\rightarrow$  Wave Funct

(everything!) = particle + wave  $\Rightarrow$  wave Equat  $\rightarrow$  wave Funct

Best we can do for either is Probability

$\int |\Psi^* \Psi| dx = \text{prob} = |\Psi|^2 dx$

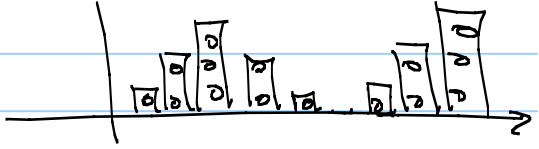
double  
slit

Ideas from Griffiths Intro to Quantum Mechanics

So... probability PART II

Keep in mind

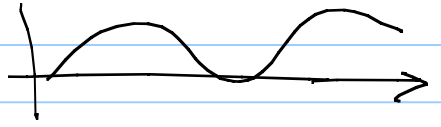
1) discrete prob =



discrete sized bins  $\Delta x$

$P_i$  for each bin

2) continuous prob  $\Rightarrow$



Lim  $\Delta x \rightarrow dx$

$\int$  prob =  $\int P(x) dx$

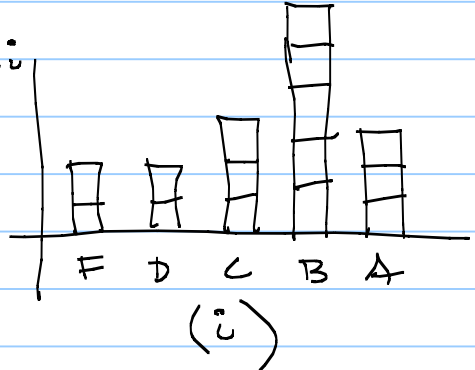
$P(x) = \text{prob density} = \frac{\text{prob}}{\text{length}}$

So  $p(x) = P(x) dx$

Start again w/ discrete prob.

TEST	$N_i$	$P_i$
A	3	$3/15$
B	5	$5/15$
C	3	$3/15$
D	2	$2/15$
F	2	$2/15$

Histogram



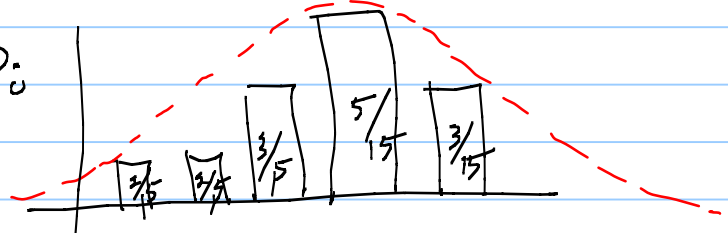
$N_{tot} = 15$

Turn to Probability (ie Normalize)

Note  $P_i = \frac{N_i}{N_{tot}}$

$P_i$

prob =



Note:  $\sum_i P_i = \frac{N_A + N_B + N_C + N_D + N_F}{15} = \frac{15}{15} = 1$  just as it should be!

Probability Histogram

add up all prob = 1

What can you do w/ your prob histogram?

Once you have  $P_i$ 's (ie prob distribution sum)

1.) Most Likely Grade ---- 1 time measurement  
 $= P_i \text{ max } (T_B = 5/15) = \text{Grade}_{\text{No. L.}}$

2.) Median  $\equiv \bar{i}$  such that  $\frac{N_{\text{tot}}}{2}$  are  $>$  &  $<$

here  $C^+/B^- = \text{Grade}_{\text{med}}$   
 $\downarrow$  new name

3.) AVERAGE = MEAN =  $\langle \text{EXPECTATION Value} \rangle$

which IS

NOT

note  $\langle \rangle$  bracket  
 $= \overline{\text{grade}} = \text{grade}_{\text{ave}} = \langle \text{grade} \rangle = \langle g \rangle$

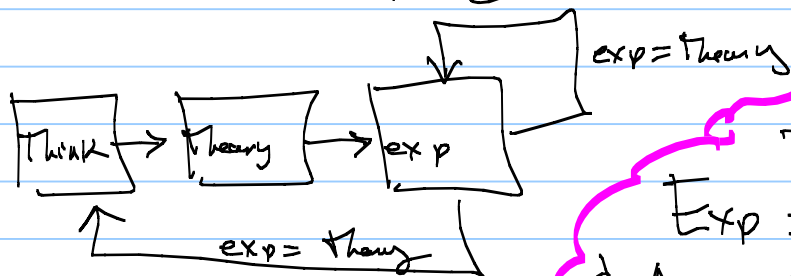
4.) What is standard deviation  
 (ie average distance from the

$$\text{Ave} = \text{mean} = \langle \text{exp} \rangle \equiv \sqrt{\sum_i \frac{(\text{grade}_i - \overline{\text{grade}})^2}{N_{\text{tot}}}}$$

We'll have to understand all 4

Why?

Scient Method



We need to know --

Exp  $\equiv$  theory

& Account for Errors that can't be part

1.) most likely -- easy,  $P_i \text{ max} = 5/15 = B = 3$

2.) median ...  $1/2$  above,  $1/2$  below  $\Rightarrow C^+ / B^- = 2.9 - 3.1$

OK now work...

3.) average = mean =  $\langle \text{expectation} \rangle$

We know

$$\overline{\text{average}} = \frac{\sum (N_i) g_i}{N_{\text{TOT}}}$$

GPA G

$$g_i: \begin{aligned} g_A &= 4 \\ g_B &= 3 \\ g_C &= 2 \\ g_D &= 1 \\ g_F &= 0 \end{aligned}$$

$$= \frac{(N_A)g_A + (N_B)g_B + (N_C)g_C + (N_D)g_D + (N_F)g_F}{N_{\text{TOT}}}$$

$$= g_A \left( \frac{N_A}{N_T} \right) + g_B \left( \frac{N_B}{N_T} \right) + g_C \left( \frac{N_C}{N_T} \right) + g_D \left( \frac{N_D}{N_T} \right) + g_F \left( \frac{N_F}{N_T} \right)$$

recognize  $\frac{N_i}{N_T} = P_i$

$$= g_A P_A + g_B P_B + g_C P_C + g_D P_D + g_F P_F$$

$$\langle \text{ave} \rangle = \sum_i g_i P_i = \text{sum over all GPA} \left( \left( \text{grade}_i \right) \left( \text{prob of grade}_i \right) \right)$$

Actually a Big result.

In general say you want the average GPA

$$\langle g \rangle = \sum_i g_i P_i$$

But then say you want over GPA<sup>2</sup>;  $\langle g^2 \rangle = \sum_i g_i^2 P_i$

So it is generally true

$$\langle f(i) \rangle = \sum_i f(i) P_i$$

ave, exp  
mean  
of  $f(i)$

Big Result

All expectations  
this way!

That's it!

Now recall most likely  $\neq$  average!

H.W. Compute for our simple examples

For our  
grade stats  $\rightarrow$

most likely

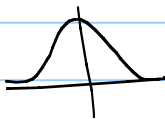
$\frac{1}{2} \langle gpa \rangle$  based on gpa

$g_A=4$   
 $g_B=3$   
 $g_C=2$   
 $SD=1$   
 $\neq \Rightarrow$

$$\langle g \rangle = \sum_i g_i P_i = g_A P_A + g_B P_B + g_C P_C + g_D P_D + g_F P_F$$

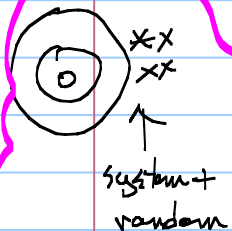
Now Bit of measurement Theory.....

Errors  $\Rightarrow$  ① part of nature - can't be completely eliminated

② Random  equal high } about true  
equal low }

tailing off far from true

= Gaussian or normal



$\neq$  Accurate (true)

$\Rightarrow$  precise (reproducible)

③ Systematic  $\Rightarrow$  part of technique of equipment that is always wrong every time.

Because we need to compare

Theory  $\stackrel{?}{\equiv}$  Experiment

we need to know

Theory  $\pm$  Theoretical error  $\stackrel{?}{\equiv}$  Experiment  $\pm$  Experiment Error

If you do not give idea you are not telling the truth & you are not doing Science!

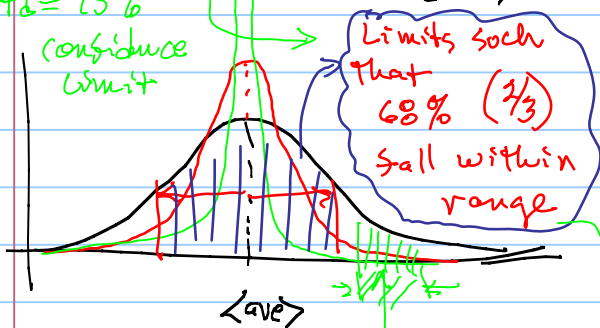
Theory  $\pm$  Theory Error

Exp  $\pm$  Exp Error

take PH414

$\langle ave \rangle \pm$  Stand deviation

measure of average distance from  $\langle ave \rangle$   
 $1 \text{ std} = 68\%$  confidence limit  
 $2 \text{ std} = 95\%$  confidence limit



1)  $\langle ave \rangle \pm \frac{\text{Stand deviat}}{\sqrt{N}}$

$N = \#$  of measurement measurement experiments

2) Counts  $\pm \sqrt{\text{Counts}}$

counting experiments

$\approx \frac{1}{3}$  = correct theory but out of this range

1) Normal distributions Gaussians

2) Binomial distributions

3 Theories but 3 different probability distributions but same  $\langle ave \rangle$

what distinguishes the 3 theories are the standard Deviations!

**NOTE it**

Experiment falls here... rules out green theory!

So... can't say enough about computing standard deviation ( $\rightarrow$  not in Scherrer --- see Griffiths)

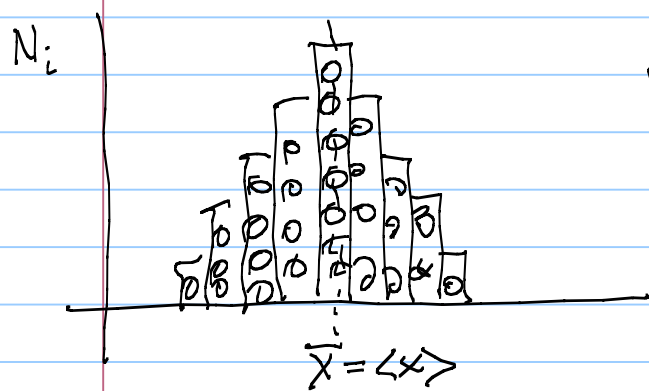
Lets try it... For  $N_{TOT}$  measurements of position  $X_i$

$$std = \sqrt{\frac{\sum (X_i - \bar{X})^2}{N_{TOT}}} \quad \text{where } \Delta \equiv |s\rangle - |i\rangle \text{ measure from } \bar{X}$$

def  $\sigma = \text{variance} = std^2$

$$\sigma = \frac{\sum (X_i - \bar{X})^2}{N_{TOT}}$$

Keep in mind orig Histogram



#'s of  $X_i$  = # in each Bin!

so

$$\sigma = \frac{\sum N_i (X_i - \bar{X})^2}{N_{TOT}}$$

$$= \sum \left[ \frac{N_i}{N_T} (X_i^2 - 2X_i\bar{X} + \bar{X}^2) \right]$$

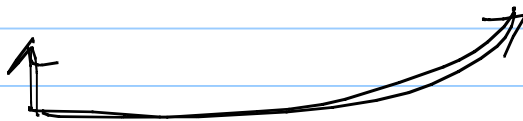
$$= \sum P_i (X_i^2 - 2X_i\bar{X} + \bar{X}^2)$$

=

$$\sigma = \sum X_i^2 P_i - 2\bar{X} \sum X_i P_i + \bar{X}^2 \sum P_i$$

Hey just def of  $\langle X \rangle = \sum X_i P_i$

$$\sigma = \langle X^2 \rangle - 2\bar{X}\bar{X} + \bar{X}^2 = \langle X^2 \rangle - \langle X \rangle^2$$



This is  $\sigma^2$ !

This is Big result!

$$\sigma_x^2 = \langle x^2 \rangle - \langle x \rangle^2$$

(expectation of  $x^2$ ) - (expectation of  $x$ )<sup>2</sup>

= variance!

std =  $\sqrt{\sigma}$  = error in theory!

1 standard deviation = 68%,  $\frac{2}{3}$  of chance answer is in this range!

Note, still

$\frac{1}{3}$  chance the answer is Good & not in this Range!

So Theory Best answer is

$$\langle x \rangle \pm \sqrt{\sigma_x}$$

→ 1std = 68% confidence range  
2std = 95% confidence range

↳ note 5% chance Theory still good w/ # is outside of  $\langle x \rangle \pm 2\sqrt{\sigma}$

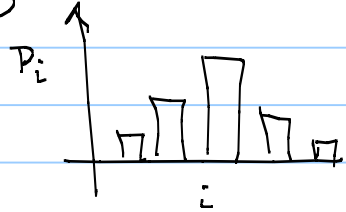


H.W. compute  $\langle g \rangle \pm \sqrt{\sigma_g}$

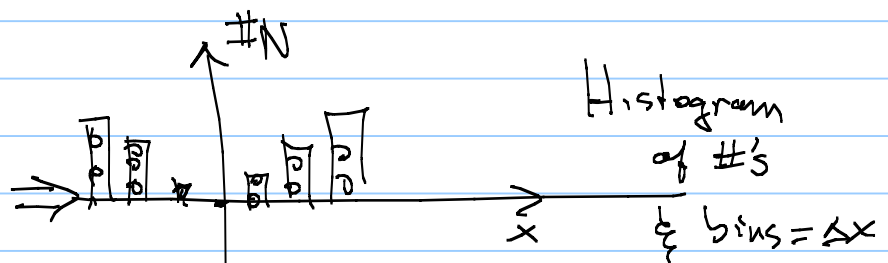
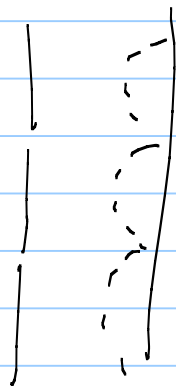


OK:

So we've got most likely  
median  
 $\langle x \rangle$   
 $\sigma$  } set discrete prob.



We've agreed that w/



but instead of labeling discrete  
bins  $\downarrow$   
ith

we let  $\Delta x \rightarrow dx$

then as  $dx \rightarrow 0$ , bin width  $\rightarrow 0$

$\downarrow$   
prob(x)  $\propto dx$   
 $\uparrow$  bin size

so

$$\text{prob}(x) = f(x) dx$$

$\uparrow$   
prob =  $\frac{\text{prob}}{\text{length}}$   
density

So... important connection is



such that

$$\langle \text{exp} \rangle = \int \Psi^* (\text{quantity}(x)) \Psi dx$$

$$= \int \text{quantity}(x) \Psi^*(x) \Psi(x) dx$$

note

if eigen problem

$$\hat{O} \Psi = o \Psi$$

But if not

$$\hat{O} \Psi = o \Psi'$$

still OK

$$\Psi^* \Psi' \text{ still} = \frac{\text{prob}}{\text{length}}$$

so

$$= \int \text{quantity}(x) \frac{\text{prob}}{\text{length}} dx$$

$$\langle \text{exp} \rangle = \int (\text{quantity}(x)) \text{prob}(x) dx$$

continuous sum over continuous values & continuous probabilities!

which looks just like discrete prob again

$$= \sum \langle \psi | \hat{O} | \psi \rangle P_i$$

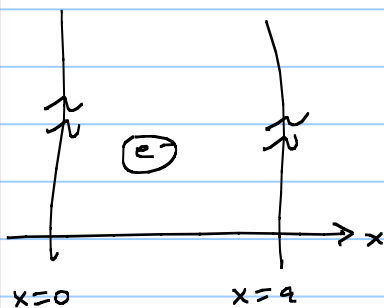
discrete sum over discrete values & probabilities

Let's make this legit...

Momentum

= Prototype  
= Sol

1) observables



$$\hat{H} \Psi(x,t) = i\hbar \frac{\partial \Psi}{\partial t}$$

$$\text{Solve } \Psi(x,t) = \begin{cases} A \sin(\frac{\pi x}{a}) e^{-i\frac{\hbar \pi^2 t}{2m a^2}} & 0 \leq x \leq a \\ 0 & \text{otherwise} \end{cases}$$

Now ask

what is average,  $\langle \text{exp} \rangle$  of momentum?

well

$$\langle p \rangle = \int \Psi^* \hat{p} \Psi dx$$

$$= \int \Psi^* (-i\hbar \frac{\partial}{\partial x}) \Psi dx$$

$$= \int \Psi^* (\hbar k) \Psi dx$$

$$= \hbar k \int \Psi^* \Psi dx$$

$$= \hbar k \underbrace{\int \Psi^* \Psi dx}_{=1} = \hbar k$$

$\hat{p} = -i\hbar \frac{\partial}{\partial x} \equiv$   
Operator  
Sol  
Linear  
momentum

\*Note: 1) Here,  $\Psi$  was eigen function of  $\hat{p}$   
But we will postulate

(just as we did in Schrod when

$$\text{added } \left( \frac{\hbar^2 k^2}{2m} + \hat{V}(x) \right) \Psi = i\hbar \frac{\partial \Psi}{\partial t}$$

that solving Sol  $\Psi$  will still work

\* we only  
justified sol  
Sol  $E_k$  not  
 $E_k + E_p$

2.) notice  $\hat{p} = \underline{\text{Linear!}}$  operator  
recall  
definition of Linear

$$\hat{O}[f(x) + g(x)] = \hat{O}f(x) + \hat{O}g(x)$$

$$\hat{O}[c f(x)] = c \hat{O}f(x)$$

3.) notice that because  
 $\Psi$  is an eigen function  
of  $\hat{p}$

you get nice definite  
result ---

$$\langle p \rangle = \hbar k$$

This is true in general

$\Psi$   
eigenfunction of  
 $\hat{O}$ 's  
represent states  
in which the  
"particle" has  
definite or "fixed"  
value for that  
observable!

→ HUGE Q.M. Discussion  
Follows next!

So when  $\bar{\Psi}$  = eigenfunction of  $\hat{O}$ , you theoretical observable is Fixed exact

$\therefore \langle P \rangle = \hbar k$  so theory has } Experiment, however will of course have error!  
 $\frac{1}{2} \langle P^2 \rangle - \langle P \rangle^2 = 0$  }  $P \pm 0$   
 $\hbar k \pm 0$

When  $\bar{\Psi} \neq$  eigenfunction of  $\hat{O}$   
 Everything holds only will get  $\langle \rangle \pm$  error!

$\implies$  when  $\bar{\Psi} \neq$  eigenfunction of  $\hat{O}$   
 will get  
 $\langle \text{expect of observable} \rangle \pm \langle \text{Error} \rangle$

> This is HUGE

Departure from classical physics: Classically you always have well defined (at least in theory)

In Q.M., unless you have eigenfunctions, you will not have definite attributes observables!

Why? ----- a Bit Later!  
 Hint ... (Heisenberg uncertainty princ.)

OK: BIG Result in Q.M.  
& Terminology from  $\hat{p} \Rightarrow$  prototype

1) Measurable quantities  $\equiv$  OBSERVABLES!  
ex: momentum = an "observable"

2) as w/  $\hat{p}$ , all observables will have  
a linear  $\hat{O}$

ex:  $\hat{D} \hat{p} = -i\hbar \frac{\partial}{\partial x}$

others  $\begin{cases} 2) \text{ position} \\ 3) \text{ Energy} \\ 4) \text{ angular momentum} \end{cases}$  we will build these 1-2, & 3

why linear?

$$\hat{L} \Psi \rightarrow c \Psi'$$

$\Psi'$  will

live in same function or vector (Hilbert) space as  $\Psi$

and  $\therefore$  Probability is preserved!

3. All observables can be built from

these (3) observable  $\hat{O}$ 's

so really only need to know these

ex: #4;  $\vec{L} = \vec{r} \times \vec{p} = \text{posit} \times \text{momentum}$

so to build 2 & 3

## Results on observables:

observable	$\hat{O}$
1) momentum: $\vec{p}$	$-i\hbar \frac{\partial}{\partial x}$ ( $-i\hbar \vec{p} = \hat{p}$ )
2) position: $\vec{x}$	$x$
3) energy: $E_{\text{TOT}}$	$i\hbar \frac{\partial}{\partial t}$

\* remember all other observable (Experimentally measurable) observables can be built from these 3

Let's do Energy next

--- we already know it!

for  $\Psi \propto e^{i(kx - \omega t)}$

we know  $E_{\text{TOT}}$  (wave-like) =  $h\nu = \frac{h}{2\pi} 2\pi\nu = \hbar\omega$

so  $\hat{E} = i\hbar \frac{\partial}{\partial t}$  does the trick!

ex:  $\langle E_{\text{TOT}} \rangle = \int \Psi^* (i\hbar \frac{\partial}{\partial t}) \Psi dx$

if  $\Psi$  = eigenfunction of  $\hat{E}$  Then

$$\langle E_{\text{TOT}} \rangle = \text{fixed value}$$

$$= \int \Psi^* \hbar\omega \Psi dx = \hbar\omega \int \Psi^* \Psi dx$$

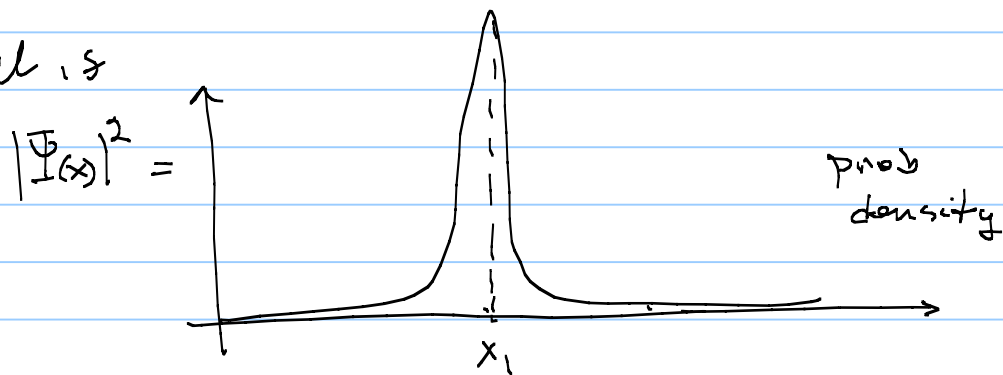
$$\langle E_{\text{TOT}} \rangle = \hbar\omega$$

Again, not all  $\Psi$  = eigenfunctions of  $\hat{E}$  so in general will get  $\langle E_{\text{TOT}} \rangle \neq \text{Error}$



and finally position  $(1-D) \hat{0}$ .

well, if



$$|\Psi(x)|^2 dx = \text{prob} = \text{max @ } x_1$$

so how to pull out  $x_1 = \text{position?}$

well

$$\int \Psi^*(x) \times x \Psi(x) dx$$

↑ doesn't do anything to  $\Psi$  as an  $\hat{0}$

$$= \int x \Psi^*(x) \Psi(x) dx$$

= prob  $\Rightarrow$  only Big near  $x_1$

$$= x_1$$

so

$$\langle x \rangle = \int \Psi^* \times x \Psi dx$$

identify

$$\langle x \rangle \text{ position } \hat{0} \approx x$$

So we've got

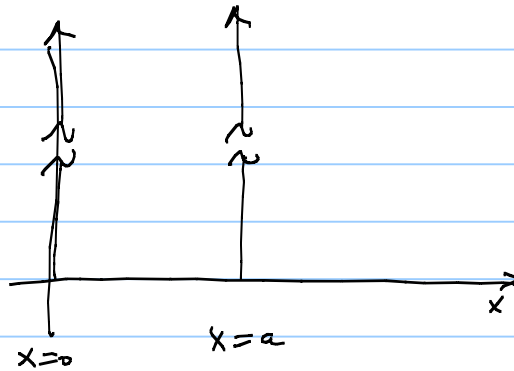
observ  $\psi$

$$x \rightarrow \hat{x}$$

$$\hat{p} \rightarrow -i\hbar \frac{\partial}{\partial x}$$

$$E_{\text{tot}} \rightarrow i\hbar \frac{\partial}{\partial t}$$

H.W. For our 1-D particle in  $\infty$  well



$$\hat{H}\Psi = i\hbar \frac{\partial \Psi}{\partial t} = \text{Schröd}$$

$$\Psi(x,t) = \begin{cases} \sqrt{\frac{2}{a}} \sin\left(\frac{n\pi x}{a}\right) e^{-i\frac{n^2 \pi^2 \hbar t}{2ma^2}} & 0 \leq x \leq a \\ 0 & \text{otherwise} \end{cases}$$

compute  $\langle x \rangle \pm \Delta x$   
 $\langle p \rangle \pm \Delta p$   
 $\langle E \rangle \pm \Delta E$

while noting, in particular,  $\Psi = \text{eigenfunc}$   
of which  $\hat{H}$  so expecting which Fixed  
observables!